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**MODELING OF LOAD TRANSFER IN CERAMIC MATRIX  
COMPOSITES**

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Esta Dissertação foi julgada adequada para obtenção do Título de Mestre em Ciência e Engenharia de Materiais, e aprovada em sua forma final pelo Programa de Pós Graduação em Ciência e Engenharia de Materiais

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All Models are wrong, but some are useful.

(George E.P. Box)





## ABSTRACT

The aim of this work is to present some models of load transfer between porous matrix and fibers in ceramic matrix composites. An analytical model for short fibers is developed, based on the earlier shear-lag models used for polymeric composites. Moreover, geometry and strength of fibers in addition to the matrix porosity are included in the present analysis. The theoretical curves for the longitudinal and shear stress distribution along the fiber-porous matrix interface are presented. They exhibited a maximum strength point at the middle of the short fibers. It became evident that the critical length is governed by the relative properties of the fibers, matrix and porosity, which greatly influenced the load carrying capacity of the fibers in the composites. In addition, the present simplified solution facilitates the understanding of the interface mechanism using porous matrix. In addition, a bundle testing routine is implemented using Monte Carlo methods. It is common knowledge that for bundles of fibers in composites, that the bundle strength is always less than the sum of the fiber strengths. This behavior can be explained by load-sharing models. At this work, different load sharing models were implemented on a simulated tensile test of ceramic oxide fibers. The results are in agreement with the experimental results of single-fiber and bundle testing and constitute a useful tool for the design of fiber-reinforced materials.

**Keywords:** modeling, load transfer, ceramic matrix composites



## RESUMO

Este trabalho se dedica a apresentar alguns modelos de transferência de carga entre uma matriz porosa e fibras em compósitos de matriz cerâmica. Um modelo analítico para a transferência de carga em fibras curtas é desenvolvido, baseado em modelos já existentes para compósitos poliméricos. Além disso, a geometria e a resistência das fibras, juntamente com a porosidade da matriz são incluídas na presente análise. As curvas teóricas para as tensões longitudinais e de cisalhamento ao longo da interface fibra-matriz são apresentadas. Elas alcançam um máximo no meio das fibras curtas. Torna-se evidente que o comprimento crítico é governado pelo conjunto de propriedades da fibra e da matriz, que influenciam a capacidade de transferência de carga nos compósitos. Adicionalmente, a solução simplificada apresentada facilita o entendimento dos mecanismos interfaciais se utilizando de uma matriz porosa. Outro foco do trabalho é um algoritmo que simula o teste de feixes contínuos de fibras cerâmicas usando-se métodos de Monte Carlo. É mostrado que a resistência do feixe é sempre menor que a resistência média das fibras testadas individualmente. Tal comportamento é explicado por modelos de transferência de carga. Neste trabalho, diferentes modelos de transferência de carga foram implementados em uma simulação de um ensaio de tração em feixes de fibras. Os resultados estão de acordo com os experimentos de fibra simples e feixe realizados e constituem uma ferramenta útil para o projeto de materiais reforçados com fibras cerâmicas.

**Palavras-chave:** modelamento, transferência de carga, compósitos de matriz cerâmica.



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## **ABBREVIATIONS AND ACRONYMS LIST**

CMCs – Ceramic Matrix Composites

SiC – Silicon Carbide

LaPO<sub>4</sub> - Monazite

ELS – Equal Load Sharing

LLS – Local Load Sharing



## SYMBOL LIST

- $F(\sigma)$  – Failure probability of a single fiber at a stress  $\sigma$   
 $G_b(\sigma)$  – Failure probability of a dry fiber bundle at a stress  $\sigma$   
 $G_i(\sigma)$  – Failure probability of an infiltrated fiber bundle at a stress  $\sigma$   
 $G_c(\sigma)$  – Failure probability of a fiber bundle within a composite at a stress  $\sigma$   
 $\alpha$  – Dunders parameter, Critical length ratio  
 $E_m$  – Matrix Young's modulus  
 $E_f$  – Fiber Young's modulus  
 $G_d$  – Energy release rate for crack deflection in an interface  
 $G_p$  – Energy release rate for crack penetration in an interface  
 $\Gamma_i$  – Interfacial toughness  
 $\Gamma_m$  – Matrix toughness (work of fracture)  
 $\Sigma$  – Crack deflection parameter  
 $L$  – Half of a fiber's length  
 $r$  – Fiber radius  
 $u$  – Displacement on the fiber  
 $v$  – Displacement on the matrix  
 $\sigma$  - Stress  
 $E$  – Young's modulus  
 $\varepsilon$  - Strain  
 $P_f$  – Load on the fiber  
 $B$  – Cox's proportionality constant  
 $m$  – Weibull Modulus  
 $\sigma_0$  – Characteristic strength ( $F=0.632$ ) for a Weibull distribution  
 $K_i$  –  $i$ -th load concentration factor for  $i$  broken neighbors  
 $k^*$  – Critical cluster size  
 $b$  – Sintering parameter  
 $v_f$  – Fiber volume fraction  
 $\sigma_f$  – Fiber rupture stress  
 $\sigma_m$  – Matrix rupture stress  
 $\sigma_T$  – Composite transversal strength  
 $\sigma_L$  – Composite longitudinal strength



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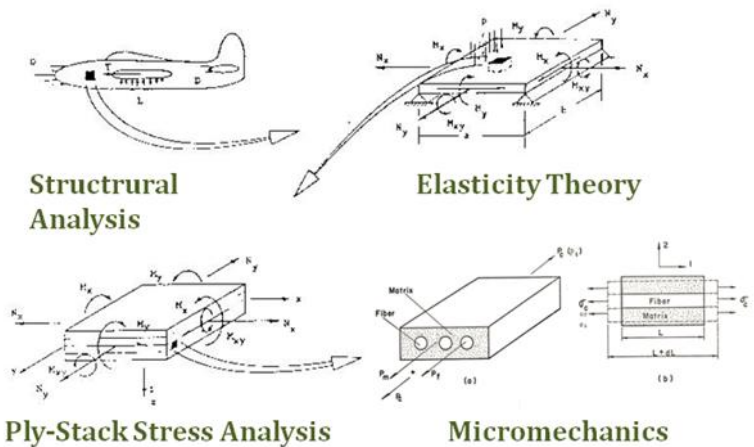
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## 1 INTRODUCTION

Modern structural ceramic composites possess a number of unique properties that cannot be achieved by other materials. Therefore, they have a potential for saving energy, reducing wear, and increasing the lifetime of components [1].

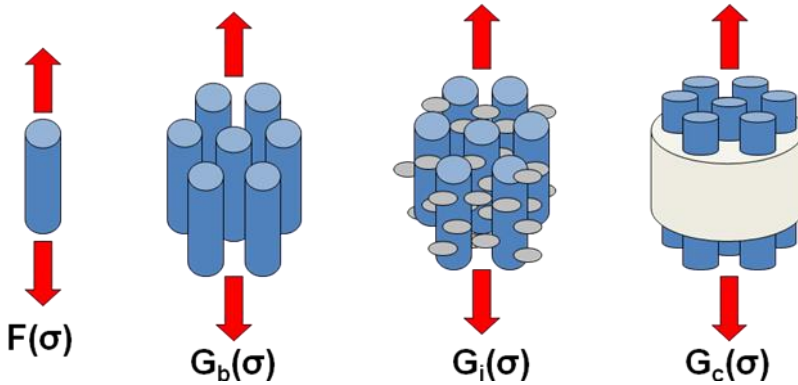
Ceramic Matrix Composites (CMCs) have attracted attention for thermomechanical applications, due to their damage tolerant fracture behavior. This is the result of toughening mechanisms, particularly crack deflection into fiber-matrix interface, as well as subsequent fiber pullout and bridging [2, 3]. Among the different categories of CMCs, all-oxide systems have recently been in the focus of research [4-9] because of their inherent high oxidation resistance compared to their non-oxide counterparts. This is interesting particularly at high temperature applications in oxidizing environments such as gas turbines.

Due to the complexity and responsibility of these materials, there is a growing need to models which can predict the bulk properties of the composite based on their microconstituents, e.g. fiber and matrix properties. This leads to micromechanical modeling (Fig. 1.1), which is an idealization of the interaction of the fibers and the matrix on the microscale.



**Fig. 1.1** Top-bottom approach for micromechanical modeling. Adapted from: [10] and [11]

The philosophy of this thesis is based on the recognition that mechanism-based models are needed, which allow for an efficient correlation to a well-conceived experimental procedure. The emphasis here is on the creation of a framework which allows models to be inserted in different complexity levels (Fig. 1.2), as they are developed, and which can also be validated by carefully chosen experiments.



**Fig. 1.2** Different complexity levels on a continuous fiber composite, each corresponding a failure probability function. From left to right: Single Fiber, Dry Bundle, Infiltrated Bundle and Consolidated Composite.

## 1.1 OBJECTIVES

This work has the main objective of understanding and modeling the mechanical behavior of ceramic matrix composites and fiber bundles, and the influence of processing and the matrix material in the mechanical behavior of the referred materials.

To achieve this goal, the following objectives were set:

- Develop a simplified shear-lag model for short-fiber ceramic matrix composites;
- Relate single-fiber properties and bundle properties;
- Simulate diverse load-sharing models for fiber bundles and determine the best suited for the studied ceramic fibers;



## 2 LITERATURE REVIEW

### 2.1 CERAMIC FIBERS

The high potential of CMCs is directly related to the use of high-resistance ceramic fibers of small diameters (usually around 10  $\mu\text{m}$ ). Covalent non-oxide fibers, as carbon or silicon carbide, are the ones showing better high-temperature mechanical properties (specially in terms of creep resistance), but are highly susceptible to oxidizing environments, calling to the use of surface treatments for protection, or the use of inert atmospheres [12].

In the other end of the spectrum, oxide fibers (as alumina and mullite-alumina), by their chemical nature, show an excellent oxidation resistance, good mechanical properties at room temperature, but present issues with creep resistance even in moderate temperatures. As consequence, the carbon and SiC fibers are the most used as reinforcement in commercial high-temperature CMCs [13].

By their small diameters, those ceramic fibers are extremely fragile and should be put into a ceramic matrix (either oxide or non-oxide), in a manner to protect them and permit the load transfer between the matrix and the fibers. The high cost of these composites is related to the high cost of those fibers, which are used in volumetric fractions ranging from 40% to 50%. Nanometric reinforcements, as carbon nanotubes, SiC nanofibers or whiskers, are not used in CMCs due to processing difficulties, cost and health hazards [12].

#### 2.1.1 Oxide Ceramic Fibers

Nextel 610 and 720 are denominations amongst a group of aluminum oxide fibers specifically designed for use as reinforcement in ceramic and metal matrix composites. Both continuous fibers are designed as composite reinforcements, but their compositional differences result in differing properties. Nextel 610 was designed to have higher strength characteristics but is susceptible to creep at elevated temperatures. Nextel 720 was then designed to have better creep resistance for elevated temperature applications, but was reduced in strength. The Nextel fibers are mostly comprised of alumina, produced via sol-gel processing, which in turn makes them less expensive to produce than some other fibers, such as SiC.

The high strength of Nextel fibers is one of its primary characteristics that make it appealing as reinforcement for composites. Their high strength is attributed in part to the fine grain structure of the material that is achieved through careful control of the processing technique. Nextel 610 fibers are comprised almost entirely of a pure  $\alpha$ -Alumina, and the Nextel 720 possess mullite specially placed on the grain boundaries. Through proper use of nucleation agents and careful control during processing, Nextel fibers are produced with a uniform microstructure comprised of grains 0.1  $\mu\text{m}$  in size and little residual porosity [12].

**Table 2.1.** Nextel 610 and 720 fiber properties [12].

<b>Property</b>	<b>Nextel 610</b>	<b>Nextel 720</b>
<b>Composition</b>	Alumina	Alumina + Mullite
<b>Weibull Modulus (m)</b>	11.4	8
<b>Characteristic Strength (MPa)</b>	3200	2200
<b>Mean Diameter (<math>\mu\text{m}</math>)</b>	10	10

## 2.2 MECHANICAL PROPERTIES

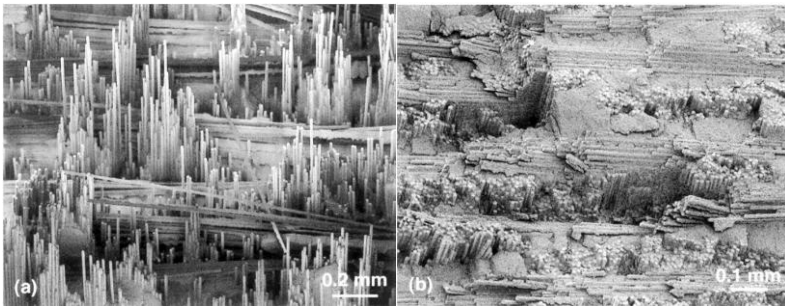
The mechanical properties of ceramic matrix composites have not been studied until the 1990's [2, 14-18]. Extensive reviews of mechanisms and mechanical properties of ceramic matrix composites are found in the literature [19, 20]. The main topics studied are dense and porous matrix composites.

For a dense-matrix composite (porosity higher than 90%), a surface treatment on the fibers is needed for crack deflection [21]. The development of oxide-oxide composites is based in a fragile fiber/matrix interface for crack deflection, giving place to oxidation-resistant coatings which are chemically stable. Monazite ( $\text{LaPO}_4$ ), hibonite and scheelite are among the various materials studied. Morgan et al. [22] and Chawla et al. [23] have shown that due to the chemical compatibility of monazite with alumina at high temperatures, this coating would be a good candidate for an interface material in alumina-based composites. Since that time, numerous manufacturing trials of monazite films and its use with different combinations of matrix and fibers were investigated

[24, 25]. The degradation of fiber resistance caused by the film and the need for expensive thermal treatments were identified as barriers to the application of these materials [25].

It was shown also that a similar behavior in relation to crack deflection can be achieved by the means of a finely distributed porosity in the matrix instead of a separate interface between matrix and fibers [14].

For a highly porous matrix, the main objective is to insulate the fibers from cracks that can start on the matrix. Due to the highly porous matrix material, the energy is dissipated and the stress concentration around the fibers is reduced. The crack propagation for the neighboring fibers is inhibited and the same are intact even with the matrix fracture (Fig. 2.1) [26].



**Fig. 2.1** Fracture surface of ceramic composites, showing: a) high-porosity matrix and b) low porosity matrix [14].

Although the matrix rules the pullout and crack deflection phenomena, the mechanical properties of the composite are strongly dependant of the fibers used as reinforcement. For composites with a volumetric fraction between 0.35 and 0.4, the typical values are of an elasticity modulus between 60 and 110 GPa and a bending strength between 140 and 220 MPa [14]. The higher values are from alumina fiber-reinforced composites (Nextel 610) and the lower from alumina-mullite fibers (Nextel 720).

This porous matrices are usually produced by pressure infiltration of slurries (Fig. 2.2) in a perform with the fibers, followed by drying and sintering [16, 27-29].

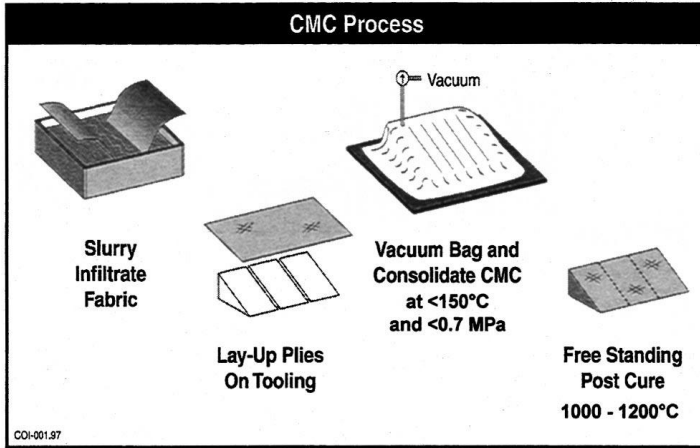


Fig 2.2 CMC production by slurry infiltration [14].

### 2.2.1 Damage Tolerance in Ceramic Composites

The damage tolerance in composite materials is thoroughly attributed to the crack deflection phenomenon between matrix and fiber (Fig. 2.3). The toughening occurs by the microcracking of the matrix and crack deflection, which keeps the fiber structure intact until the material fracture.

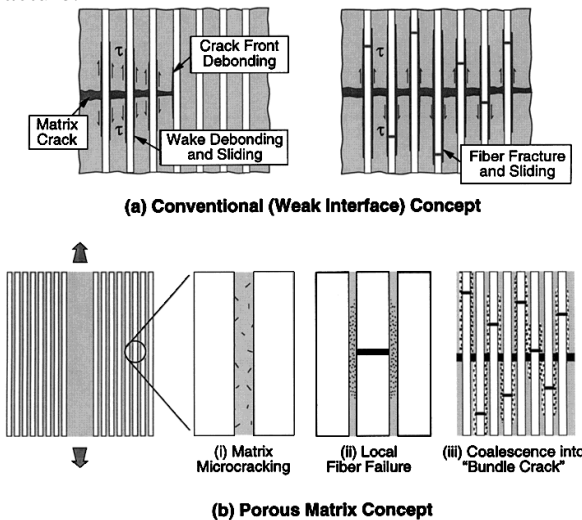


Fig. 2.3 Crack deflection phenomena in: a) dense matrix composite with weak interface and b) porous matrix composite [20].

The crack-deflection phenomena in two different materials of different elastic modulus were studied by He and Hutchinson [30]. One important variable to be considered is the Dunders parameter ( $\alpha$ ), which is a measure of the mismatch between the elastic modulus of matrix ( $E_m$ ) and fiber ( $E_f$ ):

$$\alpha = \frac{(E_f - E_m)}{(E_f + E_m)} \quad (2.1)$$

When using an energy balance, it is noted that the ratio between the energy release rate when the crack propagate between the interface  $G_d$  and the energy release rate on crack penetration  $G_p$  should be equal to the ratio of the interfacial toughness between interface and matrix [30] (eq 2.2):

$$\frac{G_d}{G_p} = \frac{\Gamma_i}{\Gamma_m} \quad (2.2)$$

A semiempiric relationship for  $G_d/G_p$  is given by Fujita et al. [27]:

$$\frac{G_d}{G_p} = \frac{1}{4(1 - \alpha)^{0.9}} \quad (2.3)$$

The graphical representation of this criterion is given by Fig. 2.4, showing where the usual porous ceramic matrix composites can be found.

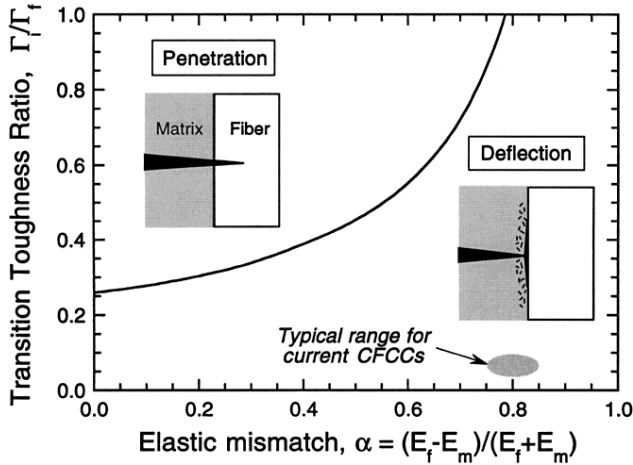


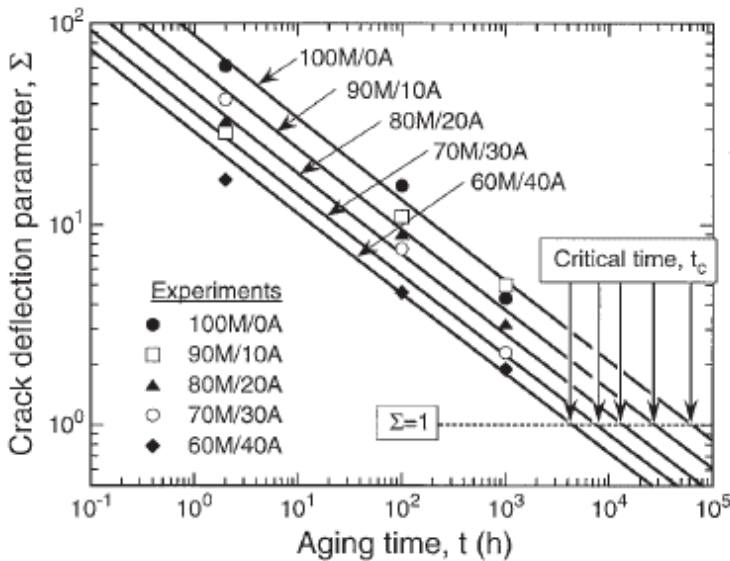
Fig. 2.4 He-Hutchinson criteria for crack deflection [20].

Replacing (2.1) in (2.3) and assuming  $\Gamma_i = \Gamma_f$ :

$$\Sigma = 0,134 \left( \frac{\Gamma_f}{\Gamma_m} \right) \left( 1 + \frac{E_f}{E_m} \right)^{0,9} \quad (2.4)$$

where  $\Sigma$  is a non-dimensional parameter which represent the propensity for crack deflection for values higher than 1. So, by knowing the relationship between the elasticity modulus between matrix and fiber, their interfacial toughness and their evolution, it is possible to predict their behavior in service and the optimal sintering parameters.

Using those criteria, Fujita et al. [27] have determined the service time of mullite-alumina composites, reinforced with Nextel 720 fibers. A model to predict the evolution of matrix properties in relation to the time was developed (Fig. 2.5):



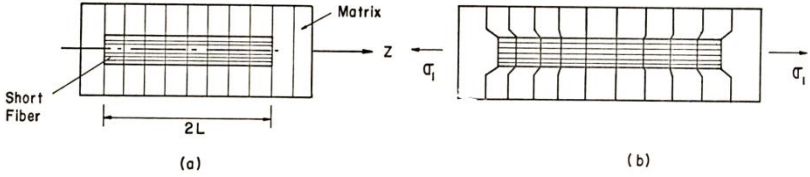
**Fig 2.5** Evolution of  $\Sigma$  with sintering time [27].

The indexes denote the matrix composition. 100M/0A would be a composite with 100% mullite and 0% alumina, and so on. Composites with a higher mullite content show a better service time, what can be explained by the lower mullite sinterability [27].

## 2.2.2 Load transfer in short fiber composites

As a pioneer model for load transfer in short-fiber reinforced composites, Cox [31] published a shear-lag model to predict the strength of paper (which is indeed a composite of cellulose and lignin fibers). The model is explained briefly in the next section:

A loaded composite made of a dense fiber with length  $2L$  is embedded in a porous matrix made of the same material as the fiber, as shown in Fig.1. It is assumed that no slippage occurs between fiber and matrix. It should also be considered that the Poisson's ratio of fiber and matrix is the same, which implies the inexistence of transversal stress when the loading is applied along the fiber. Considering the displacements in the fiber  $u$  and distant from the fiber  $v$ : (Fig. 2.6):



**Fig. 2.6** Simplified scheme of the stress field around the fiber. a) without loading. b) loaded. [11]

From Hooke's Law and taking the differential:

$$\sigma = E\epsilon = \frac{P}{A} = E \frac{\delta}{L} \quad \frac{dP}{dx} = \frac{EA}{L} \cdot \frac{d\delta}{dx} \quad (2.5)$$

Cox proposes similar behavior [31]:

$$\frac{dP_f}{dx} = B(u - v) \quad (2.6)$$

where  $P_f$  is the load acting on the fiber and  $B$  is a constant that depends on the fiber distribution and the Young's modulus of fiber and matrix.

Differentiation of Eq. 2.7 leads to:

$$\frac{d^2 P_f}{dx^2} = B \left( \frac{du}{dx} - \frac{dv}{dx} \right) \quad (2.7)$$

The derivatives of  $u$  and  $v$  can be taken as the deformations in the fiber and matrix, respectively:

$$\frac{du}{dx} = \frac{P_f}{A_f E_f} \quad (2.8)$$

$$\frac{dv}{dx} = \epsilon \quad (2.9)$$

Substitution of (2.8) and (2.9) in (2.7), gives:

$$\frac{d^2 P_f}{dx^2} = B \left( \frac{P_f}{A_f E_f} - \varepsilon \right) \quad (2.10)$$

A solution to this differential equation leads to:

$$P_f = E_f A_f \varepsilon + S \sinh(\beta x) + T \cosh(\beta x) \quad (2.11)$$

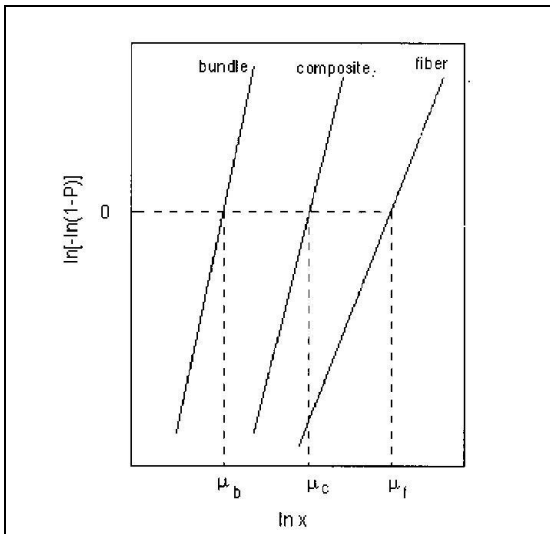
where:

$$\beta = \sqrt{\frac{B}{E_f A_f}} \quad (2.12)$$

and S and T are constants depending on the boundary conditions of the system.

### 2.2.3 Strength statistics for fiber bundles

It is well-known for bundles of fibers, that the bundle strength is always less than the sum of the fiber strengths, sometimes as much as 50% [32-37]. This is because the fibers are real materials and thus they have variable properties, and so the statistical variation needs to be taken into effect, and also the grouping and overloading effect due to the grouping. In Fig. 2.7 typical Weibull plots for single fiber strength, the strength of a bundle of these fibers, and the strength of a composite made with the bundle are shown.



**Fig. 2.7** Weibull plots for fiber tensile strength, bundle strength, and composite bundle strength [32].

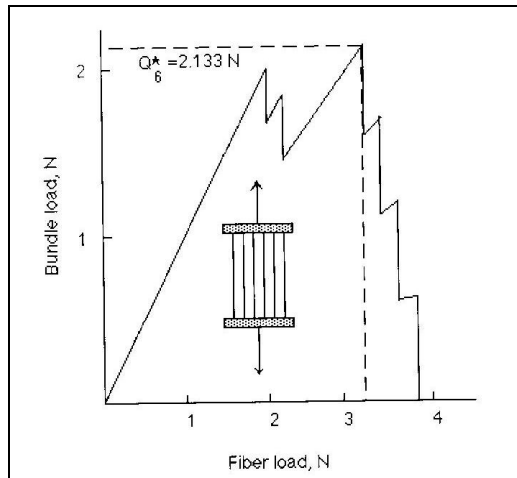


Note that going from the fiber to the bundle, the average strength is decreased, but, as the bundle is made into a composite, the strength goes up; also notice that the Weibull modulus ( $m$ ) increases, meaning the variability decreases. There are clearly things happening in the bundle and composite that cannot be explained deterministically.

### 2.2.3.1 Statistics for bundle strength. Daniel's Theorem

Consider a simple tensile experiment on a bundle of six fibers. Suppose that they are all the same size, and we know their breaking loads  $P_1 = 2.0$  N,  $P_2 = 2.2$  N,  $P_3 = 3.2$  N,  $P_4 = 3.4$  N,  $P_5 = 3.6$  N and  $P_6 = 3.8$  N. Assume that the bundle load when the load in each surviving fiber is  $P$ ,  $G_6(P)$  and denote the bundle strength by  $G_6^*(P)$ . In a deterministic world, an ultimate bundle strength  $G_6^*(P) = 3.03$ , the average fiber strength, would be the value used [32].

Then, by putting the bundle of 6 fibers in a commercial testing machine and monotonically increasing the strain, a result as the Figure 2.8 is obtained. The load in each of the fibers is identical and increases until each fiber carries a load of 2 N, and then fiber #1 fails. The surviving fibers still carry a 2 N load, but now the bundle strength is only 0.83 of its original value at the instant that the fiber broke. Now continue the extension until #2 breaks at a fiber load of 2.2 N, and the bundle strength drops again [32].



**Fig. 2.8** Simple experiment for bundle and single fiber strength [32].

Continuing on until the remaining fibers break, the peak load is found to occur when fiber #3 breaks, and this is the bundle strength  $G^*$ . A general expression for the bundle strength of a bundle with  $n$  fibers can be written [33]:

$$G_n^* = \max_{1 \leq i \leq n} \left\{ P_1, \frac{n-1}{n} P_2, \dots, \frac{n-i+1}{n} P_i, \dots, \frac{1}{n} P_n \right\} \quad (2.13)$$

More desirable, however, is being able to predict the bundle strength distribution from a knowledge of the fiber strength distribution, as well as being able to predict the strength of a large bundle of fibers; as  $n$  reaches infinity, the calculation of the former expression becomes extremely tedious. Looking more closely at equation (2.13) it can be seen that the first of the two terms is the fraction of surviving fibers while the second is the load at which they are still surviving. Motivated by this, if  $F(\sigma)$  is the failure probability for the individual fibers in the bundle, then the bundle strength,  $G^*(x)$  can be found to be [32]:

$$G^* = \left\{ \sup x \geq 0 \mid x \cdot [1 - F(\sigma)] \right\} \quad (2.14)$$

Daniels [34] was the first one to provide an analytical result to predict the bundle strength (eq. 2.15). However, it can be seen that with an increasing number of fibers, the expression itself becomes really unfeasible to calculate.

$$G_n(x) = \sum_{i=1}^n (-1)^{i+1} \binom{n}{i} F(\sigma)^i G_{n-1} \left( \frac{n\sigma}{n-i} \right), x \geq 0 \quad (2.15)$$

### 2.2.3.2 Load Sharing

In the model above it was assumed that, in a bundle under load, when a fiber fails, its load is shared equally among the surviving fibers. Such a load sharing arrangement is called an equal load sharing (ELS) rule [32-35]. Suppose the bundle load,  $G_n$ , on  $n$  fibers at the instant before the weakest fiber breaks is  $P-\varepsilon$ , where  $\varepsilon$  is very small. At this point each fiber carries the load. When the first fiber fails at  $P$ , under ELS, each of the remaining  $n-1$  fibers must be overloaded to carry the load from the broken fiber, so each fiber immediately after the breakage will bear the load  $P/(n-1)$ . The term in brackets at eq. 2.16 is called the load concentration factor, in this case  $K_1$ . In general the  $i$ th load concentration factor,  $K_i$ , under ELS is [35]:

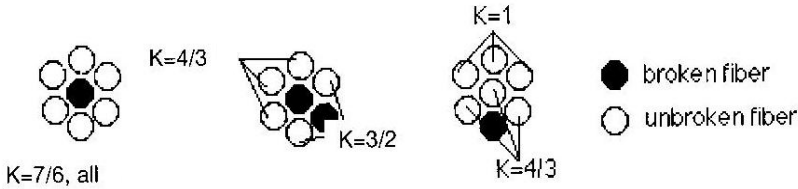
$$K_i = \left[ \frac{n}{n-i} \right], 1 \leq i \leq n-1 \quad (2.16)$$

For example, suppose a bundle has 10 fibers and the weakest fiber has strength 1. When  $G_{10} = 1$ , the first fiber will break, and immediately, each fiber will now carry a load of 1.1. At this point there are a few possibilities depending upon the strength of the next weakest fiber. If the strength is higher or equal to 1.1, all of these fibers will survive to the failure of the first. But, if only one has strength lower than 1.1, it will fail immediately and the remaining fibers will bear a 1.25 load. That means, the failure of only one fiber can lead to the catastrophic failure of the bundle relating to the overloading of the remaining fibers [32,35].

This model has some interesting implications. First, it explains why the bundle strength is lower than the mean fiber strength, as seen on Fig. 3.11. Second, it shows that the way to increase the strength of the bundle is not by simply adding stronger fibers, but rather by removing the weak ones. Because of load transfer, when many weak fibers have failed the overload will be enough to overcome any contribution of the stronger fibers. Third, another way to increase the bundle strength is that the fiber strength distribution has a high mean and as little variability as possible [32,35].

### 2.2.3.3 Local Load Sharing

The equal load sharing rule generally gives the most conservative value for bundle strength. Moreover, because the matrix in a composite tends to isolate the effects of a fiber break to the immediate vicinity of the failed fiber, the fiber's immediate and nearest neighbors bear a larger part of the overload, more than a fiber at a some distance away. A number of alternate rules to ELS have been proposed, the simplest being local load sharing, LLS [33]. Under LLS the load carried by a broken fiber is transferred only to that fiber's nearest, unbroken neighbors. Figure 2.9 illustrates this rule for several arrangements of broken fibers within a 7-fiber bundle.



**Fig. 2.9** Load intensity factors for a bundle of 7 fibers, assuming LLS [32].

Other important quantity is the number of fiber breaks required for the bundle to fail, called the critical cluster size, and is often denoted by  $k^*$ . If we know the Weibull modulus for fiber strength,  $m$ , and find the Weibull modulus for bundle strength,  $\beta$ , then the critical cluster size is [32]:

$$k^* = \frac{\beta}{m} \quad (2.17)$$

When  $k^* > 1$ , the Weibull modulus for bundle strength is higher than that for fiber strength, explaining the change in slope of the curves in Fig. 2.7.

## 3 MODELING

### 3.1 SIMPLIFIED SHEAR-LAG MODEL

#### 3.1.1 Previous Considerations and Analysis

The majority of the load transfer models for short-fiber reinforced composites was created to describe the behavior of polymer matrix composites. These include the following assumptions:

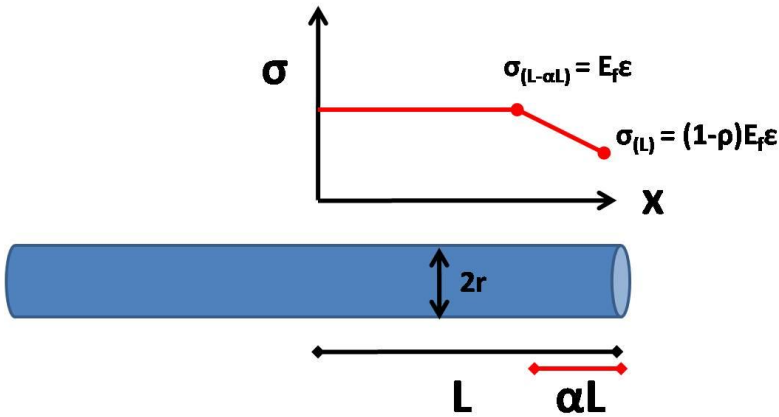
- The elastic modulus of the fibers ( $E_f$ ) is much higher than the matrix ( $E_m$ );
- The deformation until failure from the fibers ( $\epsilon_f$ ) is much lower than the matrix ( $\epsilon_m$ );
- The matrix has some degree of ductility.

Those criteria are particularly not true in the case of ceramic matrix composites, where the material of the matrix is almost the same from the fibers, so it is possible to take into account different load sharing phenomena.

The proposed model in this Thesis tries to take into account the compatibility between the fibers and matrix in porous-matrix composites, by a function of load transfer in the tip of the fibers, inversely proportional to the porosity of the matrix. Some effort is made to approximate the load transfer functions, trying to avoid the use of hyperbolic functions, which will complicate further the solution of the problem.

#### 3.1.2 Linear Shear-Lag Model

According to Fig. 3.1, let's consider a composite with fibers whose length is  $2L$ , diameter  $2r$  and Young's modulus  $E_f$ , embedded in a matrix with porosity  $\rho$ , made of the same material of the fiber. Hereby we define the critical length  $L_c$ , in which from the tip of the fiber the stress distribution isn't constant by the shear-lag between matrix and fiber. It is more feasible to work with  $\alpha$ , the ratio between the critical length and fiber length, being  $L_c = \alpha \cdot L$ .



**Fig. 3.1** Proposed stress distribution and boundary conditions

Therefore, it can be proposed that the stress distribution between the points  $L-\alpha L$  and  $L$  follows a linear behavior such as:

$$\sigma_f(x) = Ax + B \quad (3.1)$$

By using the boundary conditions defined in Fig. 3.1, and substituting then in (3.1):

$$E_f \varepsilon = A(L - \alpha L) + B \quad (3.2)$$

$$(1 - \rho)E_f \varepsilon = AL + B \quad (3.3)$$

Isolating  $B$  in (3.2) and replacing in (3.3):

$$(1 - \rho)E_f \varepsilon - AL = E_f \varepsilon - A(L - \alpha L) \quad (3.4)$$

$$-\rho E_f \varepsilon = A\alpha L \quad (3.5)$$

And then:

$$A = -\frac{\rho E_f \varepsilon}{\alpha L} \quad (3.6)$$

By replacing  $A$  from (3.1) with (3.6):

$$(1 - \rho)E_f \varepsilon = -\frac{\rho E_f \varepsilon}{\alpha} + B \quad (3.7)$$

Therefore,  $B$  is given by:

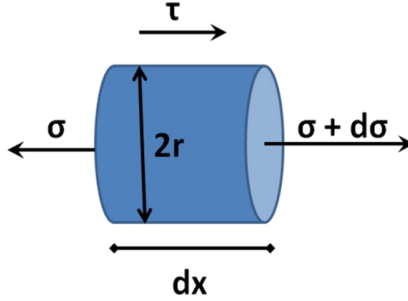
$$B = \left[ 1 - \rho \left( 1 - \frac{1}{\alpha} \right) \right] E_f \varepsilon \quad (3.8)$$

By replacing the constants in (3.1), we have the stress distribution behavior:

$$\sigma_f(x) = -\frac{\rho E_f \varepsilon}{\alpha L} x + \left[ 1 - \rho \left( 1 - \frac{1}{\alpha} \right) \right] E_f \varepsilon \quad (3.9)$$

To determine the shear stresses along the fiber, the force equilibrium in a fiber element with diameter  $2r$  and length  $dx$  is made in the  $x$  direction, resulting in:

$$\partial\sigma \cdot \pi r^2 + \tau \cdot 2\pi r dx = 0 \quad (3.10)$$



**Fig. 3.2** Force equilibrium in an infinitesimal fiber element.

Then, the shear stresses are given by:

$$\tau = -\frac{r}{2} \frac{d\sigma}{dx} \quad (3.11)$$

By the differential of (3.9):

$$\tau_f(x) = \frac{r\rho E_f \varepsilon}{2\alpha L} \quad (3.12)$$

With the stress distribution along the fiber, it is possible to calculate the average stress carried by the fiber in the composite, given by:

$$\bar{\sigma}_f = \frac{1}{L} \int_0^L \sigma_f(x) dx \quad (3.13)$$

For  $\alpha \geq 1$ , i.e. the fiber is shorter than the critical length:

$$\bar{\sigma}_f = \frac{1}{L} \int_0^L -\frac{\rho E_f \varepsilon}{\alpha L} x + \left[1 - \rho \left(1 - \frac{1}{\alpha}\right)\right] E_f \varepsilon dx \quad (3.14)$$

Then,

$$\bar{\sigma}_f = \frac{-\frac{\rho E_f \varepsilon}{2\alpha L} L^2 + \left[1 - \rho \left(1 - \frac{1}{\alpha}\right)\right] E_f \varepsilon L}{L} \quad (3.15)$$

Simplifying the equation:

$$\bar{\sigma}_f = -\frac{\rho E_f \varepsilon}{2\alpha} + \left[1 - \rho \left(1 - \frac{1}{\alpha}\right)\right] E_f \varepsilon \quad (3.16)$$

$$\bar{\sigma}_f = \left[1 - \rho \left(1 - \frac{1}{\alpha}\right) - \frac{\rho}{2\alpha}\right] E_f \varepsilon \quad (3.17)$$

Therefore, the average stress carried by the fiber is given by:

$$\bar{\sigma}_f = \left[1 - \rho - \frac{\rho}{2\alpha}\right] E_f \varepsilon \quad (3.18)$$

And for  $0 < \alpha < 1$ , i.e., the fiber is longer than the critical length:

$$\bar{\sigma}_f = \frac{\int_0^{L-\alpha L} E_f \varepsilon dx + \int_{L-\alpha L}^L \sigma_f(x) dx}{L} \quad (3.19)$$

Then,

$$\bar{\sigma}_f = \frac{E_f \varepsilon (L - \alpha L) - \frac{\rho E_f \varepsilon}{\alpha L} \int_{L-\alpha L}^L x dx + \left[1 - \rho \left(1 - \frac{1}{\alpha}\right)\right] E_f \varepsilon \int_{L-\alpha L}^L dx}{L} \quad (3.20)$$

Therefore:

$$\bar{\sigma}_f = E_f \varepsilon (1 - \alpha) - \frac{\rho E_f \varepsilon}{\alpha} \left[ \frac{1}{2} - \frac{(1 - \alpha)^2}{2} \right] + \left[1 - \rho \left(1 - \frac{1}{\alpha}\right)\right] \alpha E_f \varepsilon \quad (3.21)$$

$$\bar{\sigma}_f = E_f \varepsilon (1 - \alpha) - \rho E_f \varepsilon \left[ \frac{2 + \alpha}{2} \right] + \left[1 - \rho \left(1 - \frac{1}{\alpha}\right)\right] \alpha E_f \varepsilon \quad (3.22)$$

$$\bar{\sigma}_f = E_f \varepsilon \left(1 - \alpha - \rho + \frac{\rho \alpha}{2} + \alpha - \rho \alpha + \rho\right) \quad (3.23)$$

Simplifying the equations, we get the average stress carried by the fibers longer than the critical length:

$$\bar{\sigma}_f = E_f \varepsilon \left(1 - \frac{\rho \alpha}{2}\right) \quad (3.24)$$

With the average stresses well defined, we can define the stresses in the ply longitudinal and transversal directions. When the matrix material is the same as the fiber, it is possible to write the elastic modulus of the matrix in a function of the fiber modulus:

$$E_m = E_f e^{-bp} \quad (3.25)$$

where  $b$  is a shape factor that depends on the pore shape and distribution, according to Watchman [38].

The stress on the transversal direction is equal to the matrix maximum stress, given by:

$$\sigma_T = \sigma_m = E_f \varepsilon e^{-bp} \quad (3.26)$$

The stress on the longitudinal direction is given by the average value between matrix and fiber, based on the volumetric fractions of fiber and matrix:

$$\sigma_L = \sigma_m (1 - v_f) + \bar{\sigma}_f v_f = (1 - v_f) E_f \varepsilon e^{-bp} + \bar{\sigma}_f v_f \quad (3.27)$$

Therefore for  $0 < \alpha < 1$ :

$$\sigma_L = (1 - v_f) E_f \varepsilon e^{-bp} + E_f \varepsilon \left(1 - \frac{\rho \alpha}{2}\right) v_f \quad (3.28)$$

And for  $\alpha > 1$ :

$$\sigma_L = (1 - v_f) E_f \varepsilon e^{-bp} + E_f \varepsilon \left[1 - \rho - \frac{\rho}{2\alpha}\right] v_f \quad (3.29)$$



### 3.1.3 Quadratic Shear-Lag Model

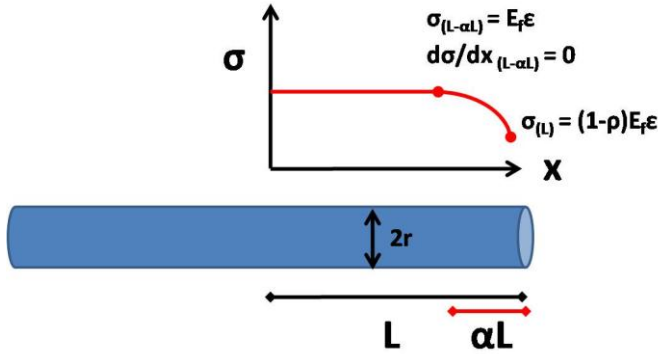


Fig. 3.3 Proposed stress distribution and boundary conditions

In a similar manner as the linear model, it can be proposed that the stress distribution between the points  $L-\alpha L$  and  $L$  follows a quadratic behavior such as:

$$\sigma_f(x) = Ax^2 + Bx + C \quad (3.30)$$

By using the boundary conditions given in Fig. 3.3, and substituting then in (3.30):

$$E_f \varepsilon = A(L - \alpha L)^2 + B(L - \alpha L) + C \quad (3.31)$$

$$(1 - \rho)E_f \varepsilon = AL^2 + BL + C \quad (3.32)$$

$$0 = 2A(L - \alpha L) + B \quad (3.33)$$

Isolating  $B$  in (3.33) and replacing in (3.31) and (3.32):

$$B = -2A(L - \alpha L) \quad (3.34)$$

$$E_f \varepsilon = A(L - \alpha L)^2 - 2A(L - \alpha L)^2 + C = -A(L - \alpha L)^2 + C \quad (3.35)$$

$$(1 - \rho)E_f \varepsilon = AL^2 - 2AL(L - \alpha L) + C \quad (3.36)$$

Subtracting (3.36) from (3.35):

$$\rho E_f \varepsilon = -A(L - \alpha L)^2 - AL^2 + 2AL(L - \alpha L) \quad (3.37)$$

$$\rho E_f \varepsilon = -AL^2(1 - 2\alpha + \alpha^2) - AL^2 + 2AL^2 - 2\alpha AL^2 \quad (3.38)$$

$$\rho E_f \varepsilon = -AL^2(1 - 2\alpha + \alpha^2 + 1 - 2 + 2\alpha) \quad (3.39)$$

$$\rho E_f \varepsilon = -A\alpha^2 L^2 \quad (3.40)$$

And then:

$$A = -\frac{\rho E_f \varepsilon}{\alpha^2 L^2} \quad (3.41)$$

By replacing  $A$  from (3.33):

$$B = -2 \left( -\frac{\rho E_f \varepsilon}{\alpha^2 L^2} \right) (L - \alpha L) \quad (3.42)$$

Therefore, B is given by:

$$B = 2 \frac{\rho E_f \varepsilon}{\alpha^2 L} (1 - \alpha) \quad (3.43)$$

To find C, we replace A in (3.34):

$$E_f \varepsilon = \frac{\rho E_f \varepsilon}{\alpha^2 L^2} (L - \alpha L)^2 + C \quad (3.44)$$

$$C = E_f \varepsilon - \frac{\rho E_f \varepsilon}{\alpha^2} (1 - \alpha)^2 \quad (3.45)$$

By replacing the constants in (3.30), we have the stress distribution behavior:

$$\sigma_f(x) = -\frac{\rho E_f \varepsilon}{\alpha^2 L^2} x^2 + 2 \frac{\rho E_f \varepsilon}{\alpha^2 L} (1 - \alpha) x + E_f \varepsilon - \frac{\rho E_f \varepsilon}{\alpha^2} (1 - \alpha)^2 \quad (3.46)$$

To determine the shear stresses along the fiber, the force equilibrium in a fiber element with diameter  $2r$  and length  $dx$  is made in the  $x$  direction, resulting in:

$$\partial \sigma \cdot \pi r^2 + \tau \cdot 2\pi r dx = 0 \quad (3.10)$$

Then, the shear stresses are given by:

$$\tau = -\frac{r}{2} \frac{d\sigma}{dx} \quad (3.11)$$

By the differential of (3.46):

$$\tau_f(x) = \frac{r \rho E_f \varepsilon}{\alpha^2 L} \left[ \frac{1}{L} x - 1 + \alpha \right] \quad (3.47)$$

With the stress distribution along the fiber, it is possible to calculate the average stress carried by the fiber in the composite, given by:

$$\bar{\sigma}_f = \frac{1}{L} \int_0^L \sigma_f(x) dx \quad (3.13)$$

For  $\alpha \geq 1$ , i.e. the fiber is shorter than the critical length:

$$\bar{\sigma}_f = \frac{1}{L} \int_0^L \left[ -\frac{\rho E_f \varepsilon}{\alpha^2 L^2} x^2 + 2 \frac{\rho E_f \varepsilon}{\alpha^2 L} (1 - \alpha) x + E_f \varepsilon - \frac{\rho E_f \varepsilon}{\alpha^2} (1 - \alpha)^2 \right] dx \quad (3.48)$$

Then,

$$\bar{\sigma}_f = \frac{-\frac{\rho E_f \varepsilon}{3\alpha^2 L^2} L^3 + \frac{\rho E_f \varepsilon}{\alpha^2 L} (1 - \alpha) L^2 + E_f \varepsilon L - \frac{\rho E_f \varepsilon}{\alpha^2} (1 - \alpha)^2 L}{L} \quad (3.49)$$

Simplifying the equation:

$$\bar{\sigma}_f = -\frac{\rho E_f \varepsilon}{3\alpha^2} + \frac{\rho E_f \varepsilon}{\alpha^2} (1 - \alpha) + E_f \varepsilon - \frac{\rho E_f \varepsilon}{\alpha^2} (1 - \alpha)^2 \quad (3.50)$$

$$\bar{\sigma}_f = \left[ 1 + \frac{\rho}{\alpha^2} \left( \frac{2}{3} - \alpha \right) - \frac{\rho}{\alpha^2} (1 - \alpha)^2 \right] E_f \varepsilon \quad (3.51)$$

$$\bar{\sigma}_f = \left[ 1 + \frac{\rho}{\alpha^2} \left( \frac{2}{3} - \alpha - 1 + 2\alpha - \alpha^2 \right) \right] E_f \varepsilon \quad (3.52)$$

Therefore, the average stress carried by the fiber is given by:

$$\bar{\sigma}_f = \left[ 1 - \rho \left( \frac{1}{3\alpha^2} - \frac{1}{\alpha} + 1 \right) \right] E_f \varepsilon \quad (3.53)$$

And for  $0 < \alpha < 1$ , i.e., the fiber is longer than the critical length:

$$\bar{\sigma}_f = \frac{\int_0^{L-\alpha L} E_f \varepsilon dx + \int_{L-\alpha L}^L \sigma_f(x) dx}{L} \quad (3.54)$$

Therefore:

$$\bar{\sigma}_f = E_f \varepsilon (1 - \alpha) - \frac{\rho E_f \varepsilon}{\alpha^2} \left[ \frac{1}{3} - \frac{(1 - \alpha)^3}{3} \right] + 2 \frac{\rho E_f \varepsilon}{\alpha^2} (1 - \alpha) \left[ \frac{1}{2} - \frac{(1 - \alpha)^2}{2} \right] \quad (3.55)$$

$$+ \left[ E_f \varepsilon - \frac{\rho E_f \varepsilon}{\alpha^2} (1 - \alpha)^2 \right] (1 - 1 + \alpha)$$

$$\frac{\bar{\sigma}_f}{E_f \varepsilon} = 1 - \alpha - \frac{\rho}{\alpha^2} \left[ \frac{1}{3} - \frac{(1 - \alpha)^3}{3} \right] + 2 \frac{\rho}{\alpha^2} (1 - \alpha) \left[ \frac{1}{2} - \frac{(1 - \alpha)^2}{2} \right] \quad (3.56)$$

$$+ \left[ 1 - \frac{\rho}{\alpha^2} (1 - \alpha)^2 \right] (1 - 1 + \alpha)$$

$$\frac{\bar{\sigma}_f}{E_f \varepsilon} = 1 - \alpha - \left[ \frac{\rho}{3\alpha^2} - \frac{\rho(1 - \alpha)^3}{3\alpha^2} \right] + 2 \frac{\rho}{\alpha^2} \left[ \frac{1}{2} - \frac{(1 - \alpha)^2}{2} - \frac{\alpha}{2} + \frac{\alpha(1 - \alpha)^2}{2} \right] \quad (3.57)$$

$$+ \left[ \alpha - \frac{\rho}{\alpha} (1 - \alpha)^2 \right]$$

$$\frac{\bar{\sigma}_f}{E_f \varepsilon} = 1 - \frac{\rho}{3\alpha^2} [1 - 1 + 3\alpha - 3\alpha^2 + \alpha^3] \quad (3.58)$$

$$+ \frac{\rho}{\alpha^2} [1 - 1 + 2\alpha - \alpha^2 - \alpha + \alpha - 2\alpha^2 + \alpha^3] - \frac{\rho}{\alpha} (1 - \alpha)^2$$

$$\frac{\bar{\sigma}_f}{E_f \varepsilon} = 1 - \frac{\rho}{\alpha} \left[ 1 - \alpha + \frac{\alpha^2}{3} \right] + \frac{\rho}{\alpha} [2 - 3\alpha + \alpha^2] - \frac{\rho}{\alpha} (1 - \alpha)^2 \quad (3.59)$$

$$\frac{\bar{\sigma}_f}{E_f \varepsilon} = 1 - \frac{\rho}{\alpha} \left[ 1 - \alpha + \frac{\alpha^2}{3} - 2 + 3\alpha - \alpha^2 + 1 - 2\alpha + \alpha^2 \right] \quad (3.60)$$

Simplifying the equations, we get the average stress carried by the fibers longer than the critical length:

$$\bar{\sigma}_f = E_f \varepsilon \left( 1 - \frac{\rho\alpha}{3} \right) \quad (3.61)$$

With the average stresses well defined, we can define the stresses in the ply longitudinal and transversal directions. When the matrix material is the same as the fiber, it is possible to write the elastic modulus of the matrix in a function of the fiber modulus:

$$E_m = E_f e^{-\rho b} \quad (3.25)$$

where  $b$  is a shape factor that depends on the pore shape and distribution, as discussed previously.

The stress on the transversal direction is equal to the matrix maximum stress, given by:

$$\sigma_T = \sigma_m = E_f \varepsilon e^{-\rho b} \quad (3.26)$$

The stress on the longitudinal direction is given by the average value between matrix and fiber, based on the volumetric fractions of fiber and matrix:

$$\sigma_L = \sigma_m(1 - v_f) + \bar{\sigma}_f v_f = (1 - v_f)E_f \varepsilon e^{-bp} + \bar{\sigma}_f v_f \quad (3.27)$$

Therefore for  $0 < \alpha < 1$ :

$$\sigma_L = (1 - v_f)E_f \varepsilon e^{-bp} + E_f \varepsilon \left(1 - \frac{\rho\alpha}{3}\right) v_f \quad (3.62)$$

And for  $\alpha > 1$ :

$$\sigma_L = (1 - v_f)E_f \varepsilon e^{-bp} + E_f \varepsilon \left[1 - \rho \left(\frac{1}{3\alpha^2} - \frac{1}{\alpha} + 1\right)\right] v_f \quad (3.63)$$

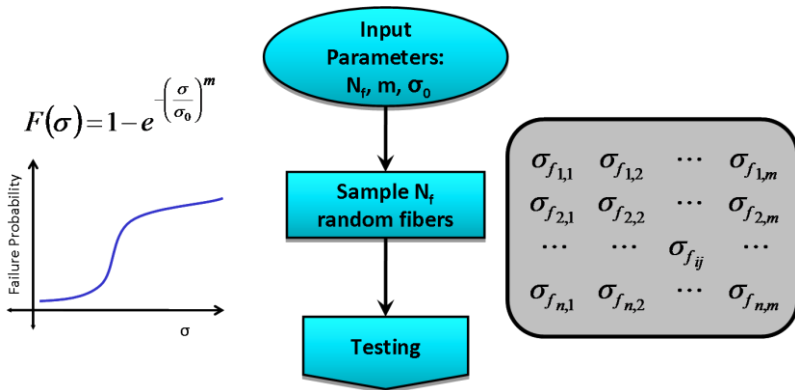
### 3.2 MONTE CARLO SIMULATION OF BUNDLE TESTING

The approach used to predict the ceramic bundle strength was a Monte-Carlo simulation of a tensile bundle test of dry fibers. The Matlab algorithm consisted of two main steps: generation of a random fiber bundle based on the Weibull parameters of single-fiber testing (Fig. 3.4) and simulated test of the created bundle (Fig. 3.5).

A Matlab routine was created in order to simulate the mechanical behavior of fiber bundles, with different load sharing rules, as a way to take into account the effects of processing and matrix in the fiber bundles.

The main steps on the simulation are the following:

- Generation of bundle of  $n$  fibers via a random fiber population from input Weibull parameters ( $m$  and  $\sigma_0$ );
- Increasing the load stepwise and individually compares it with the fibers. If the load is not enough to break a fiber, the load is increased. Otherwise, the compared fiber is broken and the load is redistributed according to the load-sharing rule;
- The above step is repeated until all fibers are broken;
- The ultimate load is recorded and the whole procedure is repeated 50 times in order to obtain a Weibull distribution;
- The program calculates the output Weibull parameters in bundle testing.



**Fig 3.4** Scheme of the bundle generation algorithm.

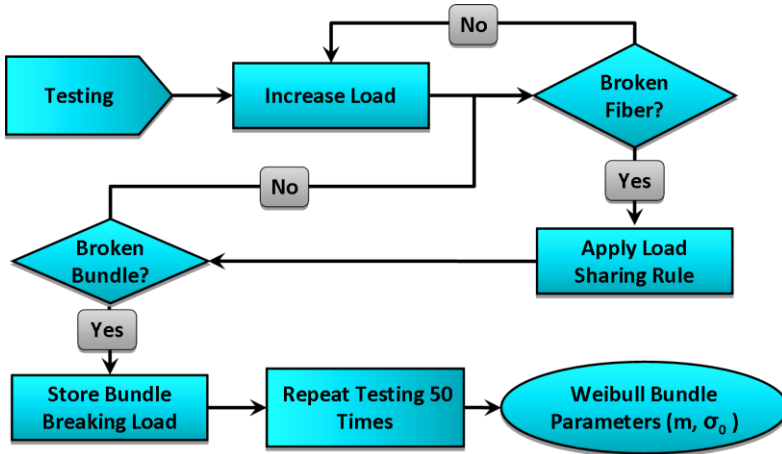


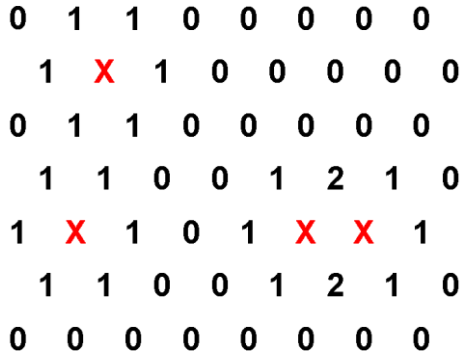
Fig 3.5 Scheme of the bundle testing algorithm.

### 3.2.1 Implementation of Load Sharing

The basis for the implementation of the load sharing is in the concept of load concentration factor,  $K$ . The bundle is seen by the program as a matrix of  $N \times M$  fibers, each with a random breaking load, based on the Weibull distribution of the single fiber data.

The program compares this bundle-matrix with the load in the machine, if one fiber breaks, this load is multiplied by a load concentration matrix,  $K$ , which has also  $N \times M$  items. In the case of equal load sharing, this factor is simply the total of fibers in the bundle divided by the number of remaining fibers.

In the case of local load sharing, whenever a fiber fails, it is marked and the program counts for each fiber the number of fractured neighbors, as can be seen in Fig. 3.6 for a hexagonal array. The failed fibers are the red Xs.



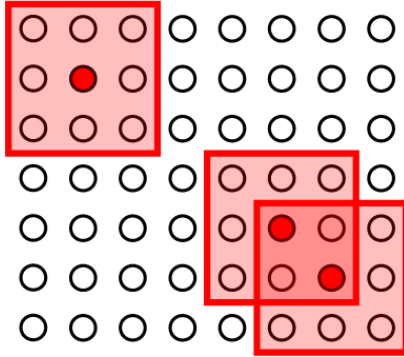
**Fig 3.6** Neighbor counting in a hexagonal array.

Then, the load concentration factor is calculated from the literature, based on the number of failed neighbors, according to Table 3.1.

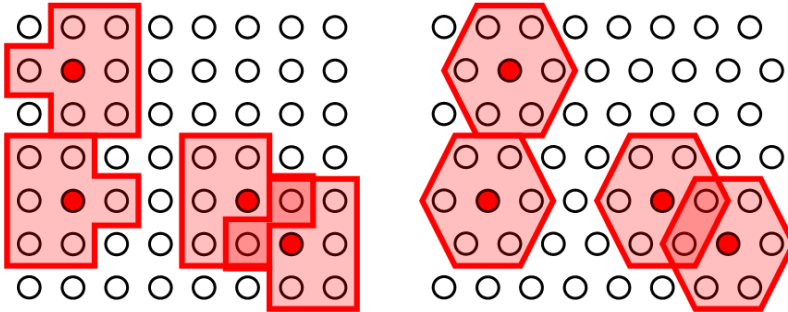
**Table 3.1** Load concentration factors

<b>Number of Broken Neighbors</b>	<b>Circular LLS Rule</b>	<b>Argon, Elastic Matrix</b>	<b>Zweben and Rosen</b>
<b>0</b>	1	1	1
<b>1</b>	1.5	1.49	1.33
<b>2</b>	2	1.76	1.6
<b>3</b>	2.5	1.92	1.83
<b>4</b>	3	2.07	2.03
<b>10</b>	6	2.72	2.97

Also, the neighbor counting method can be done in two ways: Considering a square (Fig. 3.7) or a hexagonal (Fig. 3.8) array. The implementation of the hexagonal array on a matrix is also shown, just being implemented by conditional counting in odd or even rows.



**Fig 3.7** Neighbor counting for a square array.



**Fig 3.8** Neighbor counting for a hexagonal array.

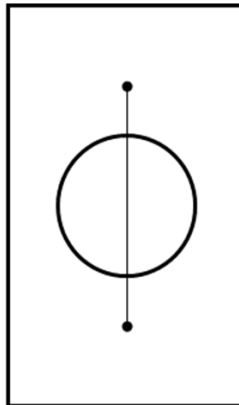


## 4 MATERIALS AND METHODS

### 4.1 FIBER PREPARATION AND SAMPLE MOUNT DESIGN

Textiles of Nextel 610 fibers were obtained from 3M for the purposes of this study. The fiber bundles were carefully separated from the textiles and the fibers were desized according to the manufacturer's recommendations. The Nextel fibers could not be easily placed into the testing grips, due to their small size and fragile nature. Through multiple trials, key aspects that came to light regarding the testing of individual fibers included fiber handling, successfully loading fibers for testing, and preserving fibers so that fracture surfaces of the tested fibers could be examined. As a result, a sample mount technique was adapted from techniques available on the literature and modified to fit with this examination [39].

Providing support for handling of the Nextel fibers, while still allowing for the ease of tensile testing, was of main importance. Index cards were cut to 70 mm in length and 50 mm in width, with a hole with a diameter of 25mm punched in the center (Fig. 4.1). A fiber would then be glued into place on the card using superglue (cyanoacrylate glue). Once secured in the tensile grips, the card was then separated into two separate pieces through the use of a scissor. The same approach was used to the tensile testing of bundles, although the literature [40] recommends different gripping methods, in order to produce comparable results between single-fiber and bundle testing.



**Fig. 4.1** Single fiber specimen mounted on the clamps for testing.

## 4.2 TENSILE TESTING

The tensile testing of single fibers and bundles (1500 den, ~400 fibers per bundle) was conducted with a controlled load on a Instron testing machine, with a 5N and 200 kN (for single-fiber and bundle tests, respectively) load cell using fiber tension test clamps. The fibers were tested using a controlled deformation mode, with preloading and a constant displacement ramp rate of 1 mm/min to a maximum of 4000 MPa. At least 29 specimens were tested in order to determine the statistical distribution.

## 4.3 DATA TREATMENT

In order to observe the statistical nature of the fiber and bundle strength, the resulting values on the mechanical testing were plotted according to Weibull's distribution (4.1).

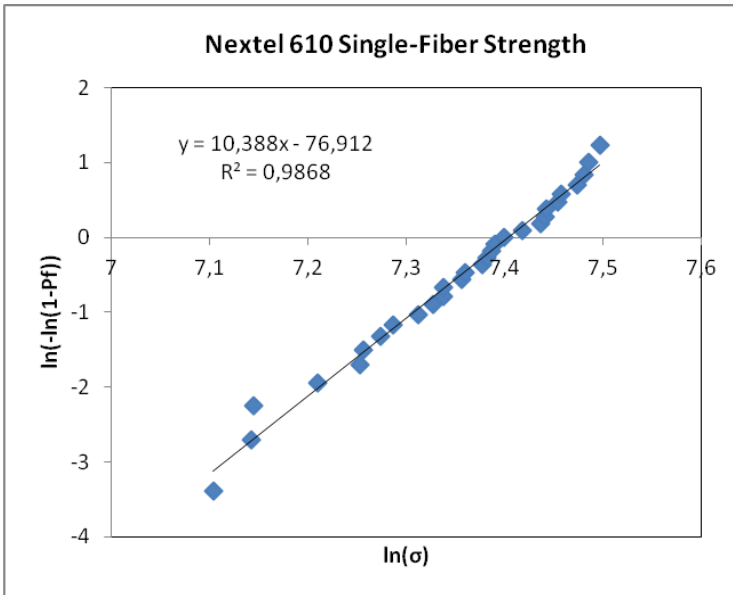
$$P_f = 1 - e^{-\left(\frac{\sigma}{\sigma_0}\right)^m} \quad (4.1)$$

The mechanical testing data was ranked and each one was given a failure probability of  $n/N+1$ , where  $n$  is the rank of the data and  $N$  is the total number of tests. Those values were fitted with the linearized form of the distribution (4.2), yielding to the  $m$  and  $\sigma_0$  values of the distribution (Fig. 4.2).

$$\ln(-\ln(1 - P_f)) = m \ln \sigma - m \ln \sigma_0 \quad (4.1)$$

**Table 4.1** Data Treatment for the fiber testing.

<i>Data Rank</i>	<i>P<sub>f</sub></i>	<i>Load (N)</i>	<i>ln(-ln(1-P<sub>f</sub>))</i>	<i>Tensile Strength (MPa)</i>	<i>ln(σ)</i>
<b>1</b>	0,033	47,3	-3,384	1215,9	7,103
<b>2</b>	0,067	49,2	-2,697	1264,7	7,142
<b>3</b>	0,1	49,3	-2,250	1267,3	7,144
<b>4</b>	0,133	52,6	-1,944	1352,1	7,209
<b>5</b>	0,166	54,9	-1,701	1411,3	7,252
...	...	...	...	...	...



**Fig. 4.2** Weibull fit of the single-fiber testing.



## 5 RESULTS AND DISCUSSION

### 5.1 SHEAR-LAG MODEL THEORETICAL RESULTS

To evaluate the models herein described, it is possible to apply the equations to an idealized composite, made of a porous alumina matrix and alumina fibers. The following table summarizes the important properties, taken as typical values from the literature:

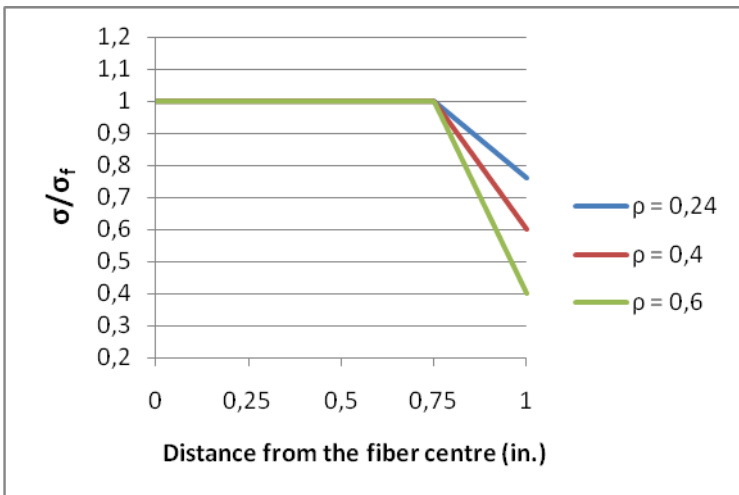
**Table 5.1** Simulated Composite Properties.

Property	Value
Fiber Volume Fraction	0.45
Matrix Porosity (%)	24
Fiber Length – 2L (mm)	50.8
Fiber Diameter (μm)	10
Critical Length / Length Ratio (α)	0.25

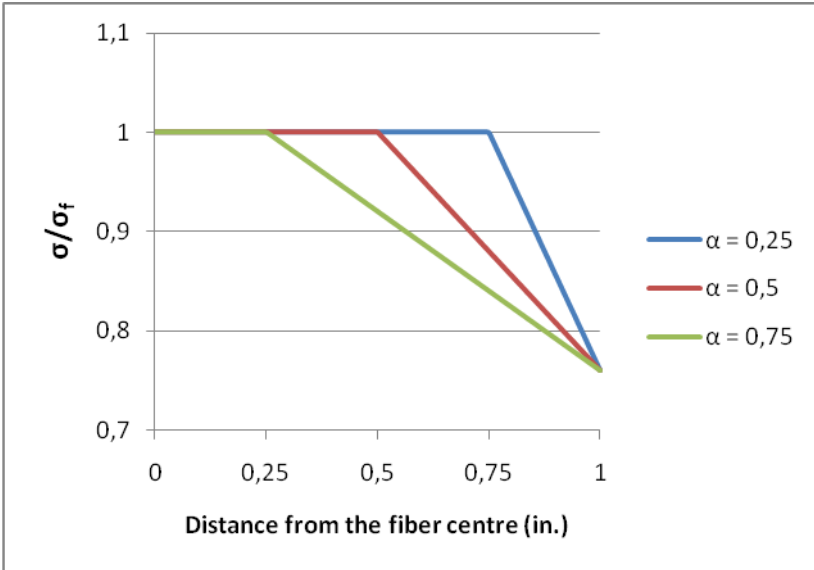
#### 5.1.1 Linear Shear-Lag Model

##### 5.1.1.1 Stress distribution

$$\sigma_f(x) = -\frac{\rho E_f \varepsilon}{\alpha L} x + \left[1 - \rho \left(1 - \frac{1}{\alpha}\right)\right] E_f \varepsilon \quad (3.9)$$



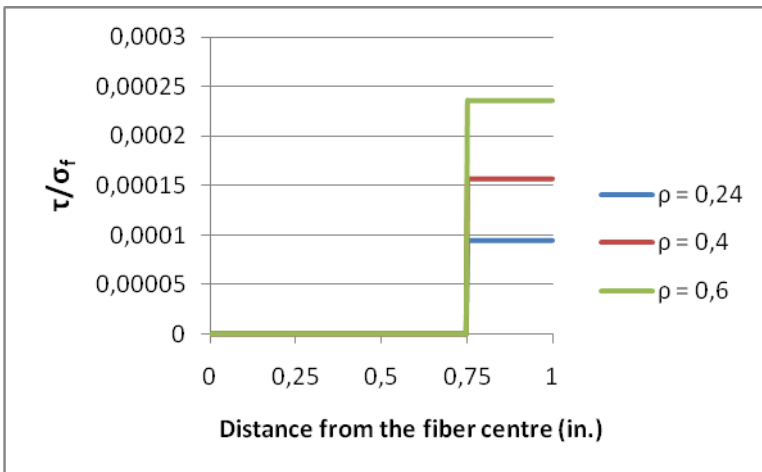
**Fig. 5.1** Stress distribution along the fiber, for different matrix porosities.



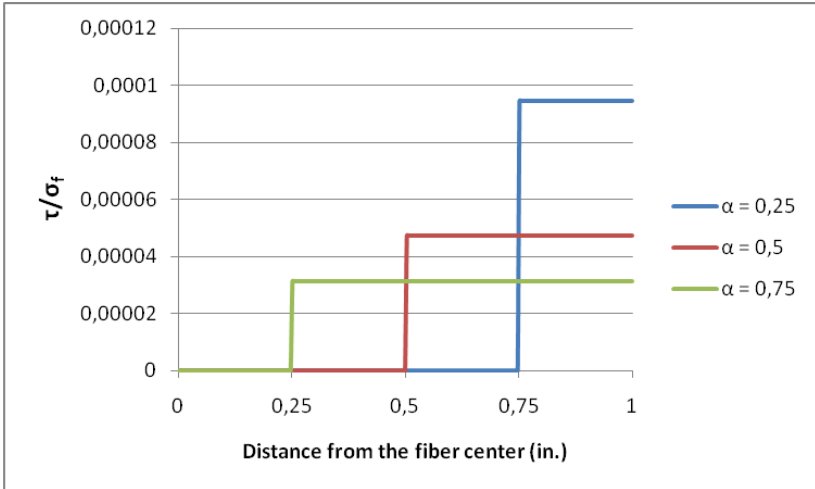
**Fig. 5.2** Stress distribution along the fiber, for critical length ratios.

### 5.1.1.2 Shear Stresses

$$\tau_f(x) = \frac{r\rho E_f \varepsilon}{2\alpha L} \quad (3.12)$$



**Fig. 5.3** Shear stress distribution along the fiber, for different matrix porosities.

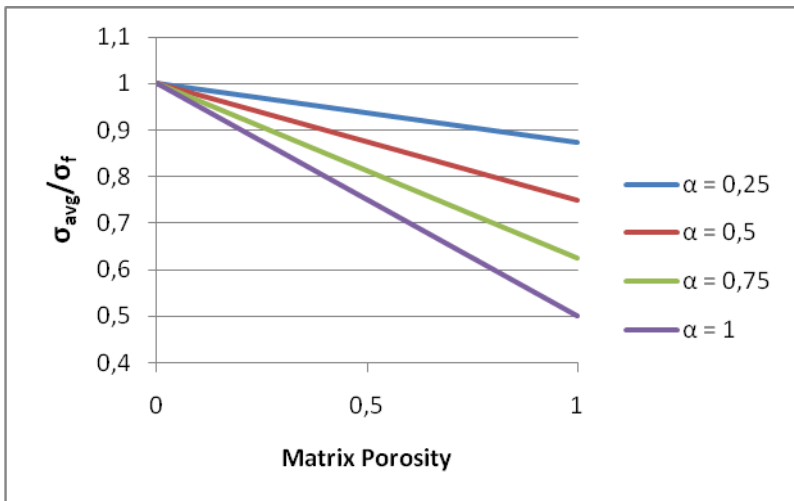


**Fig. 5.4** Shear stress distribution along the fiber, for critical length ratios.

### 5.1.1.3 Average Stresses

$$\bar{\sigma}_f = \left[ 1 - \rho - \frac{\rho}{2\alpha} \right] E_f \varepsilon \quad \text{for } \alpha > 1 \quad (3.18)$$

$$\bar{\sigma}_f = E_f \varepsilon \left( 1 - \frac{\rho\alpha}{2} \right) \quad \text{for } 1 \geq \alpha \geq 0 \quad (3.24)$$



**Fig. 5.5** Average stress carried by the fiber, for critical length ratios.

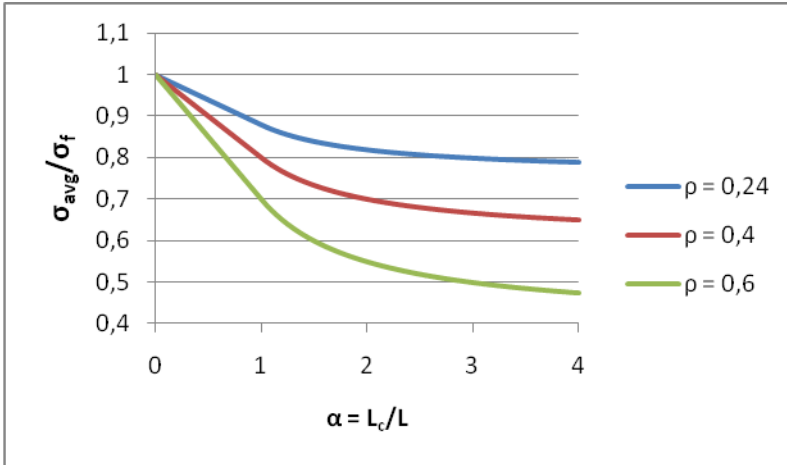


Fig. 5.6 Average stress carried by the fiber, for different matrix porosities.

#### 5.1.1.4 Longitudinal Ply Strength

$$\sigma_L = (1 - v_f)E_f \varepsilon e^{-b\rho} + E_f \varepsilon \left(1 - \frac{\rho\alpha}{2}\right) v_f \quad \text{for } 0 < \alpha < 1 \quad (3.28)$$

$$\sigma_L = (1 - v_f)E_f \varepsilon e^{-b\rho} + E_f \varepsilon \left[1 - \rho - \frac{\rho}{2\alpha}\right] v_f \quad \text{for } \alpha > 1 \quad (3.29)$$

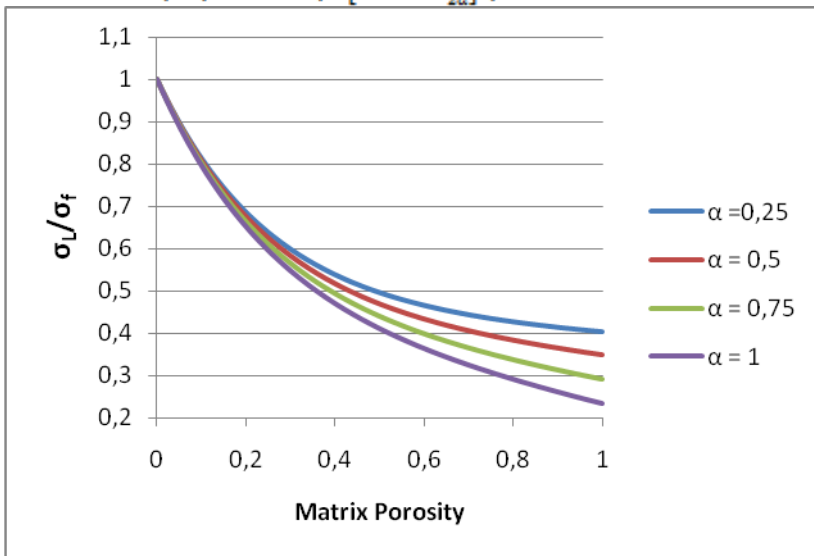
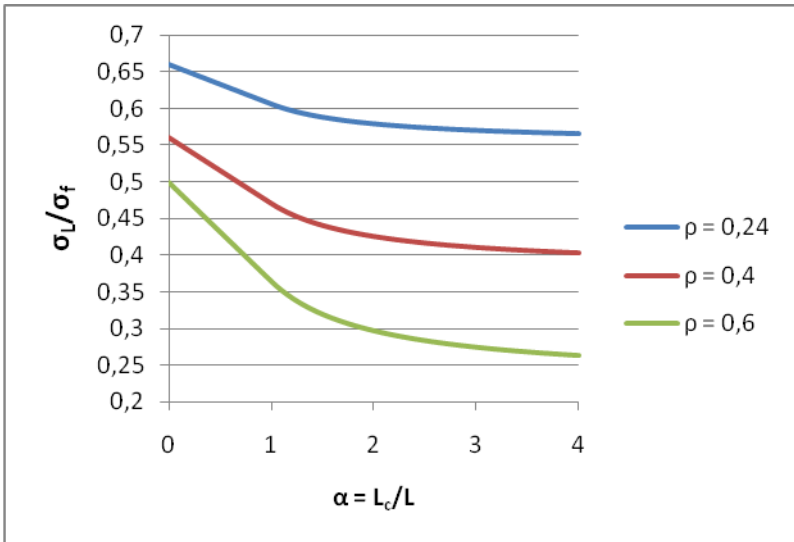


Fig. 5.7 Longitudinal Ply Strength, for critical length ratios.



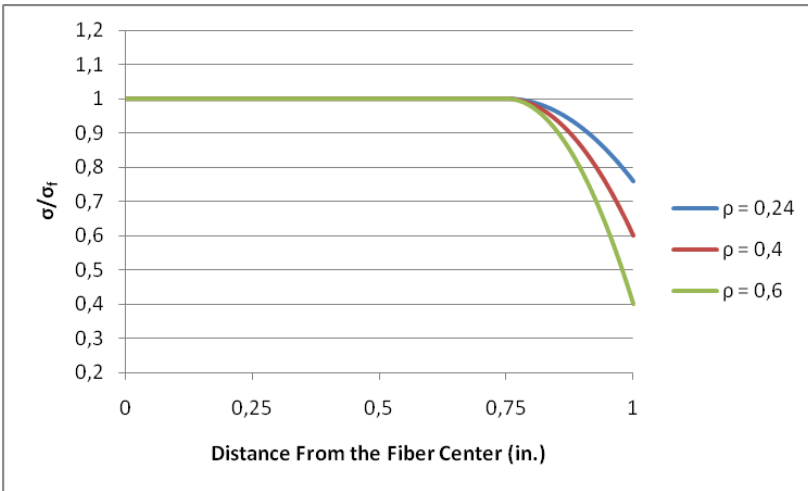


**Fig. 5.8** Longitudinal Ply Strength, for different matrix porosities.

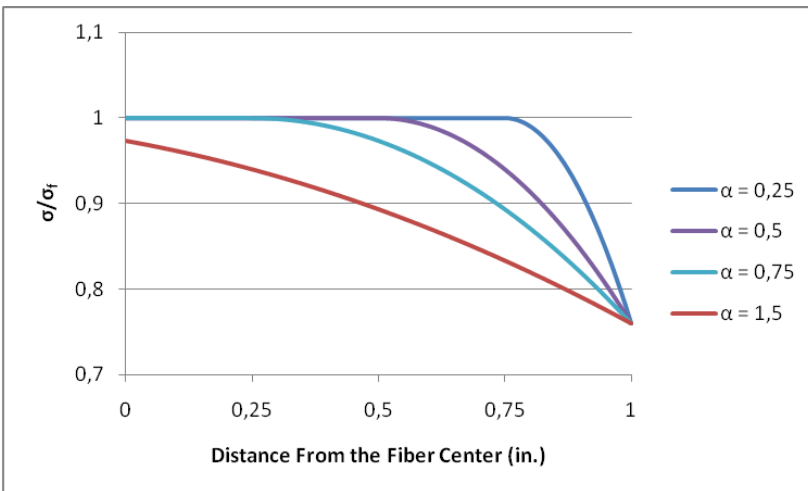
## 5.1.2 Quadratic Shear-Lag Model

### 5.1.2.1 Stress distribution

$$\sigma_f(x) = -\frac{\rho E_f \varepsilon}{\alpha^2 L^2} x^2 + 2 \frac{\rho E_f \varepsilon}{\alpha^2 L} (1 - \alpha) x + E_f \varepsilon - \frac{\rho E_f \varepsilon}{\alpha^2} (1 - \alpha)^2 \quad (3.46)$$



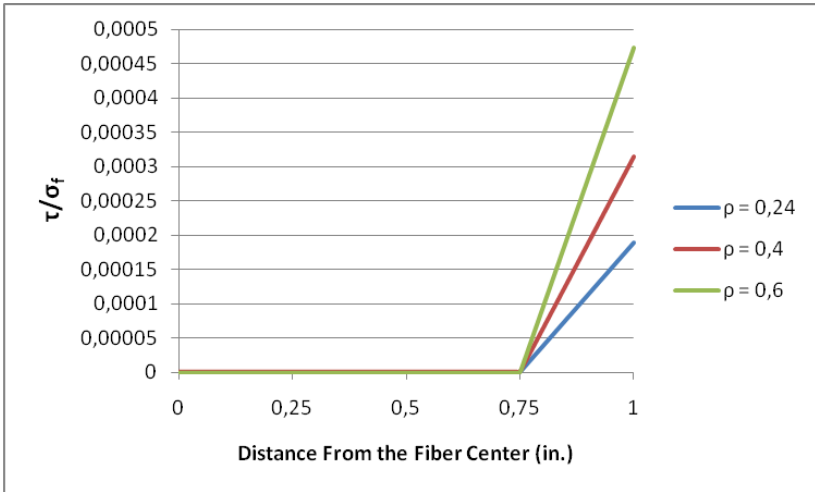
**Fig. 5.9** Stress distribution along the fiber, for different matrix porosities.



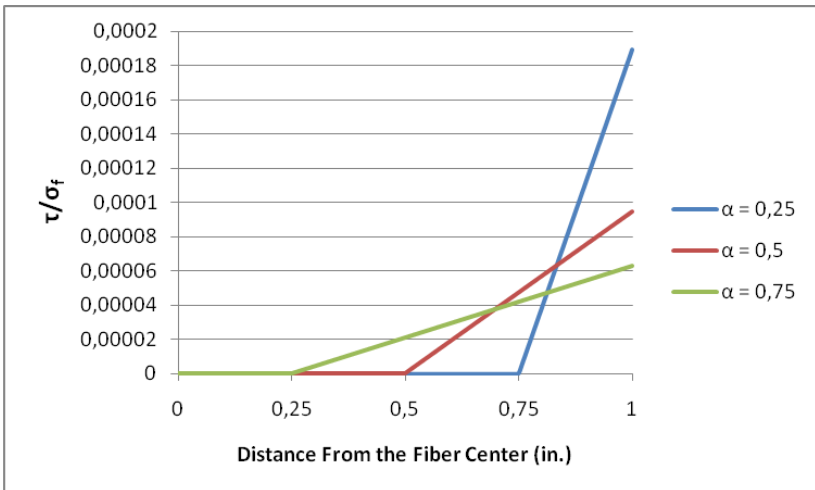
**Fig. 5.10** Stress distribution along the fiber, for critical length ratios.

## 5.1.2.2 Shear Stresses

$$\tau_f(x) = \frac{r\rho E_f \varepsilon}{\alpha^2 L} \left[ \frac{1}{L} x - 1 + \alpha \right] \quad (3.47)$$



**Fig. 5.11** Shear stress distribution along the fiber, for different matrix porosities.

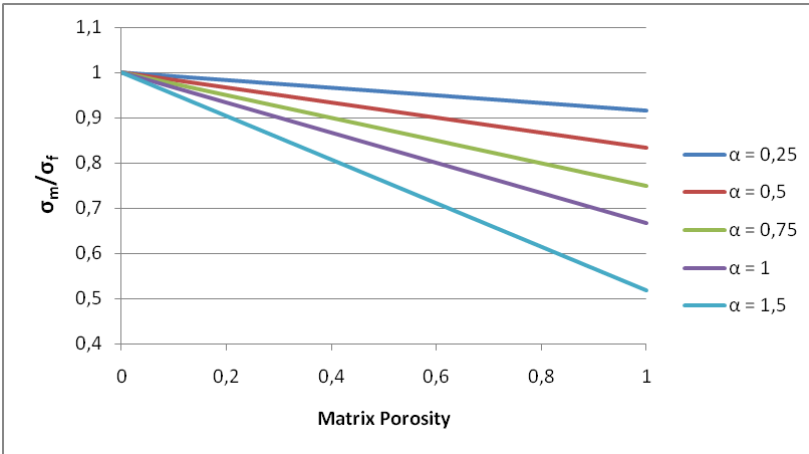


**Fig. 5.12** Shear stress distribution along the fiber, for critical length ratios.

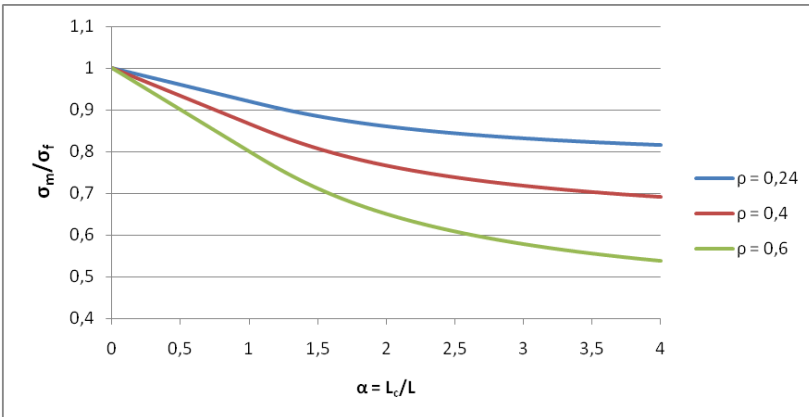
## 5.1.2.3 Average Stresses

$$\bar{\sigma}_f = \left[ 1 - \rho \left( \frac{1}{3\alpha^2} - \frac{1}{\alpha} + 1 \right) \right] E_f \varepsilon \quad \text{for } \alpha > 1 \quad (3.53)$$

$$\bar{\sigma}_f = E_f \varepsilon \left( 1 - \frac{\rho\alpha}{3} \right) \quad \text{for } 1 \geq \alpha \geq 0 \quad (3.61)$$



**Fig. 5.13** Average stress carried by the fiber, for critical length ratios.

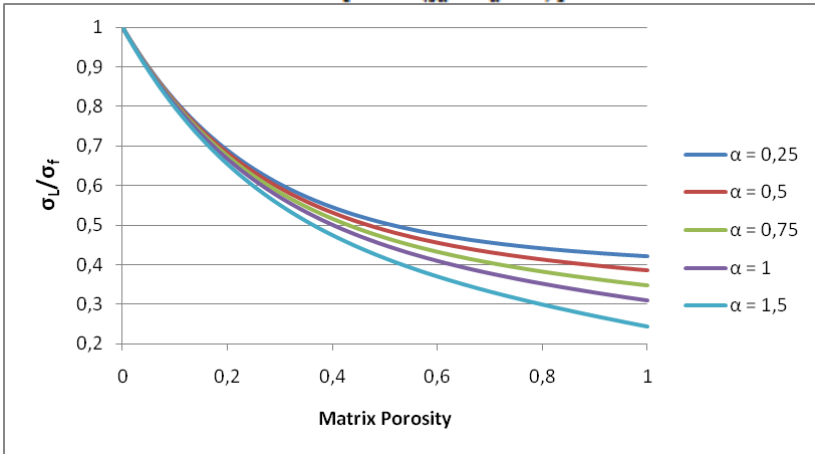


**Fig. 5.14** Average stress carried by the fiber, for different matrix porosities

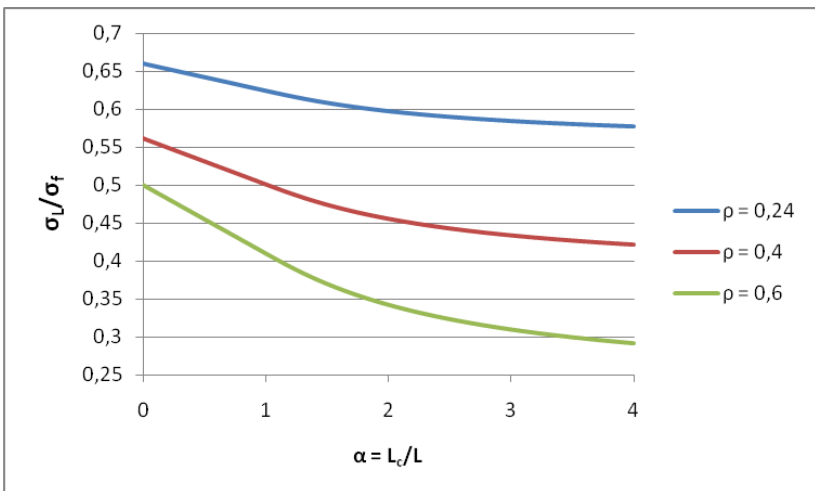
### 5.1.2.4 Longitudinal Ply Strength

$$\sigma_L = (1 - \nu_f) E_f \varepsilon e^{-b\rho} + E_f \varepsilon \left(1 - \frac{\rho\alpha}{3}\right) \nu_f \quad \text{for } 0 < \alpha < 1 \quad (3.62)$$

$$\sigma_L = (1 - \nu_f) E_f \varepsilon e^{-b\rho} + E_f \varepsilon \left[1 - \rho \left(\frac{1}{3\alpha^2} - \frac{1}{\alpha} + 1\right)\right] \nu_f \quad \text{for } \alpha > 1 \quad (3.63)$$



**Fig. 5.15** Longitudinal Ply Strength, for critical length ratios.



**Fig. 5.16** Longitudinal Ply Strength, for different matrix porosities.

### 5.1.3 Comparison with Literature

As to evaluate the effectiveness of the models developed, model predictions are compared to a porous silicon carbide matrix composite reinforced with random-aligned silicon carbide fibers, as reported by Qin et al. [41].

The following parameters are assumed in order to make the calculations:

**Table 5.2** Simulated Composite Properties.

Property	Value
<b>Fiber Volume Fraction [41]</b>	0.53
<b>Fiber Length – 2L (mm) [41]</b>	0.3-1
<b>Fiber Diameter (μm) [41]</b>	13
<b>Bulk bending strength (MPa) [41]</b>	300
<b>Critical Length / Length Ratio, α</b>	1
<b>Fiber and bulk density (g/cm<sup>3</sup>) [41]</b>	2.5
<b>Sintering parameter, b [38]</b>	4

The matrix porosity was obtained from the published composite densities, using the law of mixtures [11], leading to the following equation:

$$P_m = 1 - \frac{\rho_c - v_f \rho_f}{\rho_{th}(1 - v_f)} \quad (5.1)$$

**Table 5.3** Simulation Results.

Sintering Temperature	1650 °C	1750 °C
<b>Composite Density (g/cm<sup>3</sup>)</b>	2.03	2.46
<b>Matrix Porosity (%)</b>	40	3.4
<b>Measured Bending Strength (MPa) [41]</b>	50.75	155.75
<b>Predicted Strength (MPa) – Linear [Error]</b>	47.55 [6.31%]	168.31 [8.07%]
<b>Predicted Strength (MPa) – Quadratic [Error]</b>	50.73 [0.04%]	168.58 [8.24%]

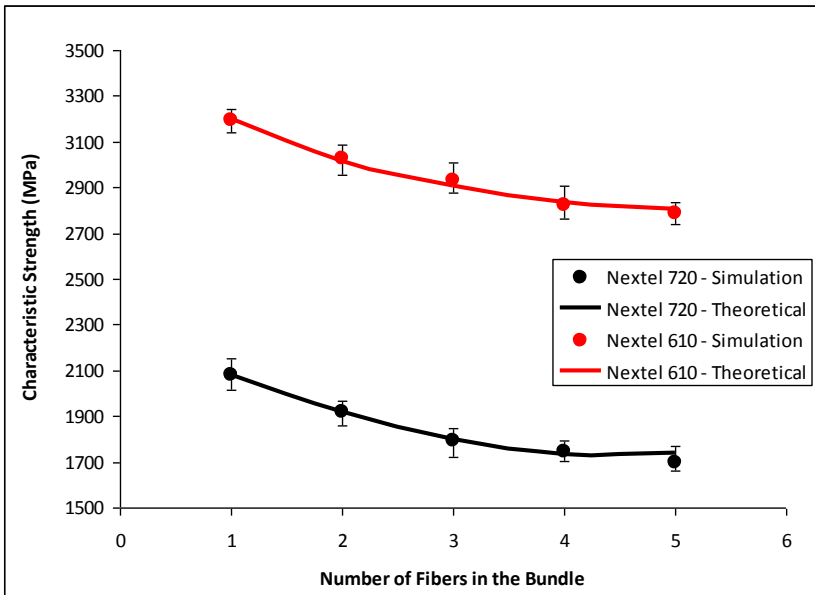
As can be seen, the predictions are in a good agreement with the experimental values reported on the literature, even with considerable simplifications leading to the calculation of matrix porosity and the determination of bulk bending strength. The difference between the linear and quadratic model predictions isn't negligible and both models provide a good range of predictions, considering the boundary conditions adopted in this case.

## 5.2 MONTE CARLO SIMULATION RESULTS

### 5.2.1 Theoretical Tests for ELS

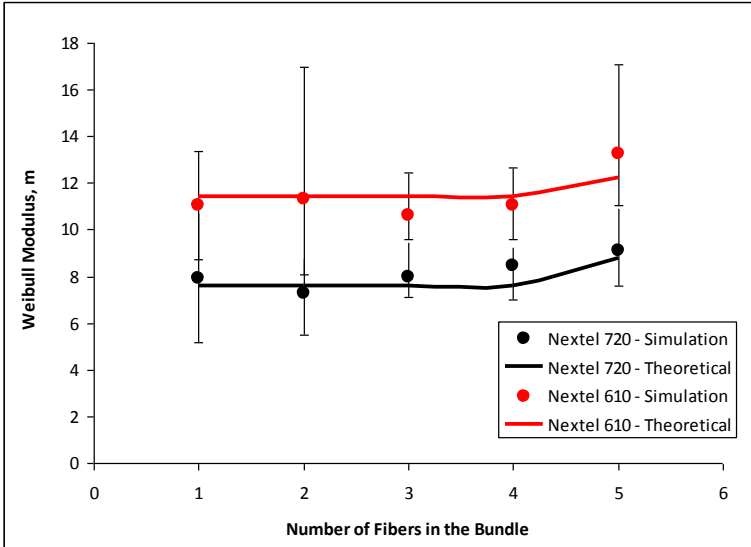
As a way to test the accuracy of the program, some tests were performed to compare its results to the analytical expressions derived by Daniels (eq 2.15).

Test runs with one to five fibers in the bundle were performed and the results were compared to the theoretical predictions based on Daniels' Theory. The fiber input data was as provided from the manufacturer, and as can be seen, both the characteristic strength ( $\sigma_0$ ) and Weibull modulus ( $m$ ) are successfully predicted in these conditions with the Equal Load Sharing algorithm.



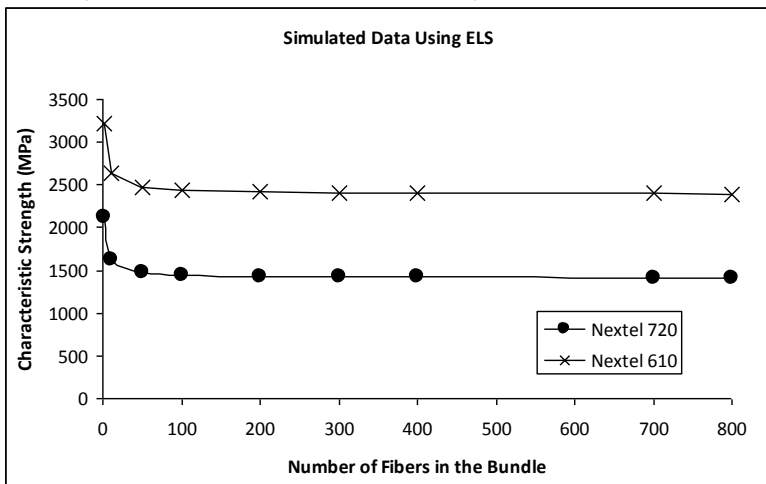
**Fig. 5.17** Simulation for ELS, dependence of characteristic strength with increasing number of fibers.



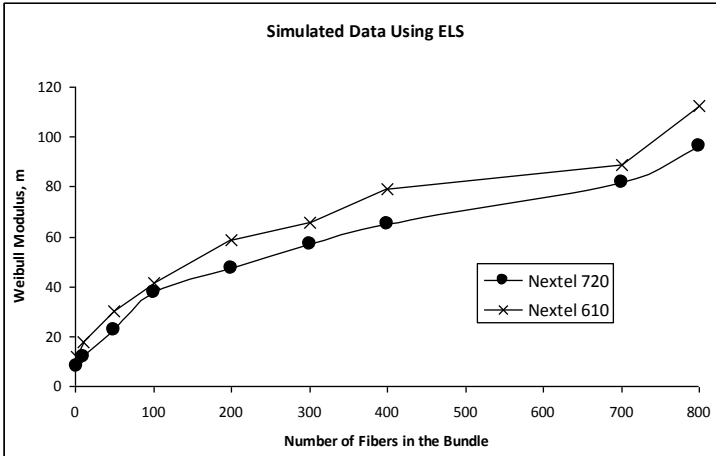


**Fig. 5.18** Simulation for ELS, dependence of Weibull modulus with increasing number of fibers.

The test runs were also made with a higher number of fibers in the simulated bundle. The results of the evolution of  $\sigma_0$  and  $m$  with increasing number of fibers are shown in Figs (5.19) and (5.20).



**Fig. 5.19** Simulation for ELS, dependence of characteristic strength with increasing number of fibers.

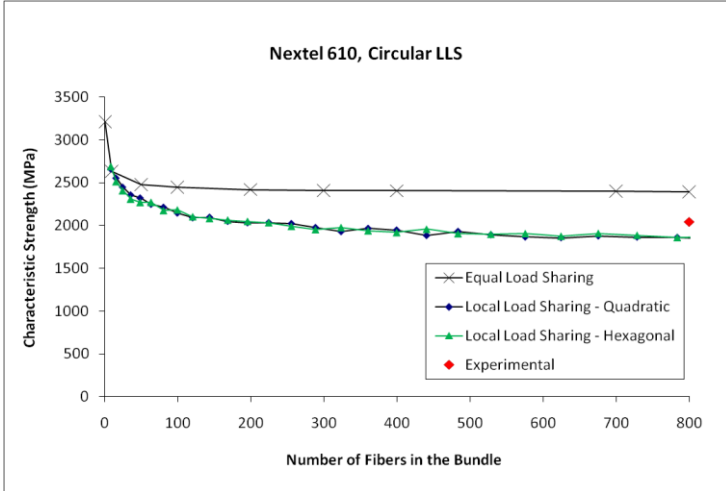


**Fig. 5.20** Simulation for ELS, dependence of Weibull modulus with increasing number of fibers.

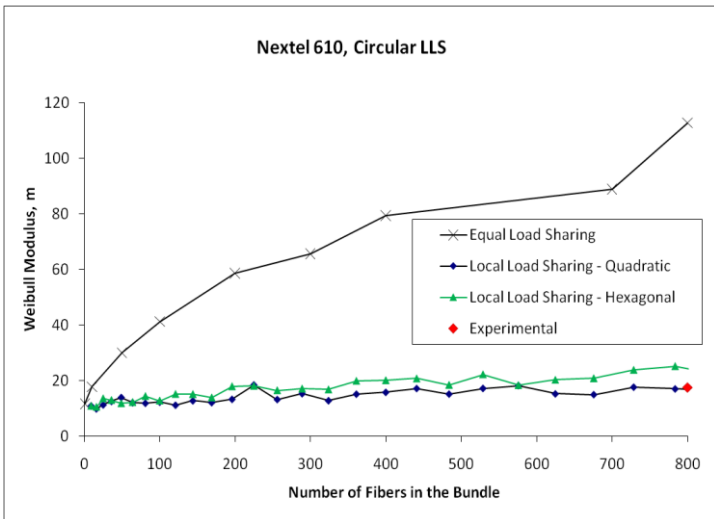
Note that for an increasing number of fibers in the bundle, the characteristic strength reaches a limit, just as like the equation (2.14), showing that the numerical routine follows the analytical reasoning. One interesting result is in Fig 5.20. It shows that under ELS, the Weibull modulus increases to unrealistic amounts. This shows clearly that even within a dry bundle, the increasing number of fibers also isolate local failures, and the theoretical prediction of ELS are unsuitable for a high number of fibers in the bundle.

## 5.2.2 Simulation Results for ELS and LLS

Figs. (5.21) and (5.22) show the evolution of the Weibull parameters in the LLS simulations using a Circular LLS rule for the stress intensity factors.



**Fig. 5.21** Simulation for LLS, dependence of characteristic strength with increasing number of fibers.



**Fig. 5.22** Simulation for ELS, dependence of Weibull modulus with increasing number of fibers.

It can be seen that the LLS theory is more suitable for a bundle with a higher number of fibers, even for dry, desized bundles. One reasonable explanation can be that with the increasing number of fibers, the slippage and friction between the fibers can transmit some part of the overloading locally via shear stress, like the bundles infiltrated with a consolidated matrix.

## 6 CONCLUDING REMARKS

This thesis developed some models of load transfer between porous matrix and fibers in ceramic matrix composites, concerning short-fiber reinforced composites with a porous matrix, and the mechanical behavior of dry fiber bundles.

An analytical model for short fibers was developed, based on the earlier shear-lag models used for polymeric composites. Moreover, geometry and strength of fibers in addition to the matrix porosity were included in the present analysis. The theoretical curves for the longitudinal and shear stresses distributions along the fiber -porous matrix interface were presented. It became evident that the critical length is governed by the relative properties of the fibers, matrix and porosity, which greatly influenced the load carrying capacity of the fibers in the composites. In addition, the present simplified solution facilitates the understanding of the interface mechanism (shear stress transfer) using porous matrix.

Using data from experiments in the literature, the model was validated, predicting in a successful manner the bending strength of SiC short-fiber reinforced silicon carbide, predicting the influence of the porosity of the matrix.

In addition, a bundle testing algorithm using Monte Carlo Methods was developed. The local-load sharing model results were in a good agreement with the experimental results of single-fiber and bundle testing, showing that for even dry fiber bundles some degree of local load sharing due to friction and slippage. Further development in the model is being made, in order to include factors as damage in the handling of the fibers and slurry infiltration. The model proved flexible and resilient enough to be further complicated.



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SILVA, J. G. P. ; HOTZA, D. ; JANSSEN, R. ; AL-QURESHI, H. A. . Modelling of load transfer between porous matrix and short fibres in ceramic matrix composites. **WIT transactions on engineering sciences (Online)**, v. 72, p. 165-174, 2011.

SILVA, J. G. P. ; AL-QURESHI, H. A. ; HOTZA, D. . Simplified Theoretical Analysis of Short Fibers in Porous Ceramic Matrix. In: 7th International Conference on High Temperature Ceramic Materials and Composites, 2010, Bayreuth. **Proceedings of the 7th International Conference on High Temperature Ceramic Materials and Composites**, 2010. p. 221-227.

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